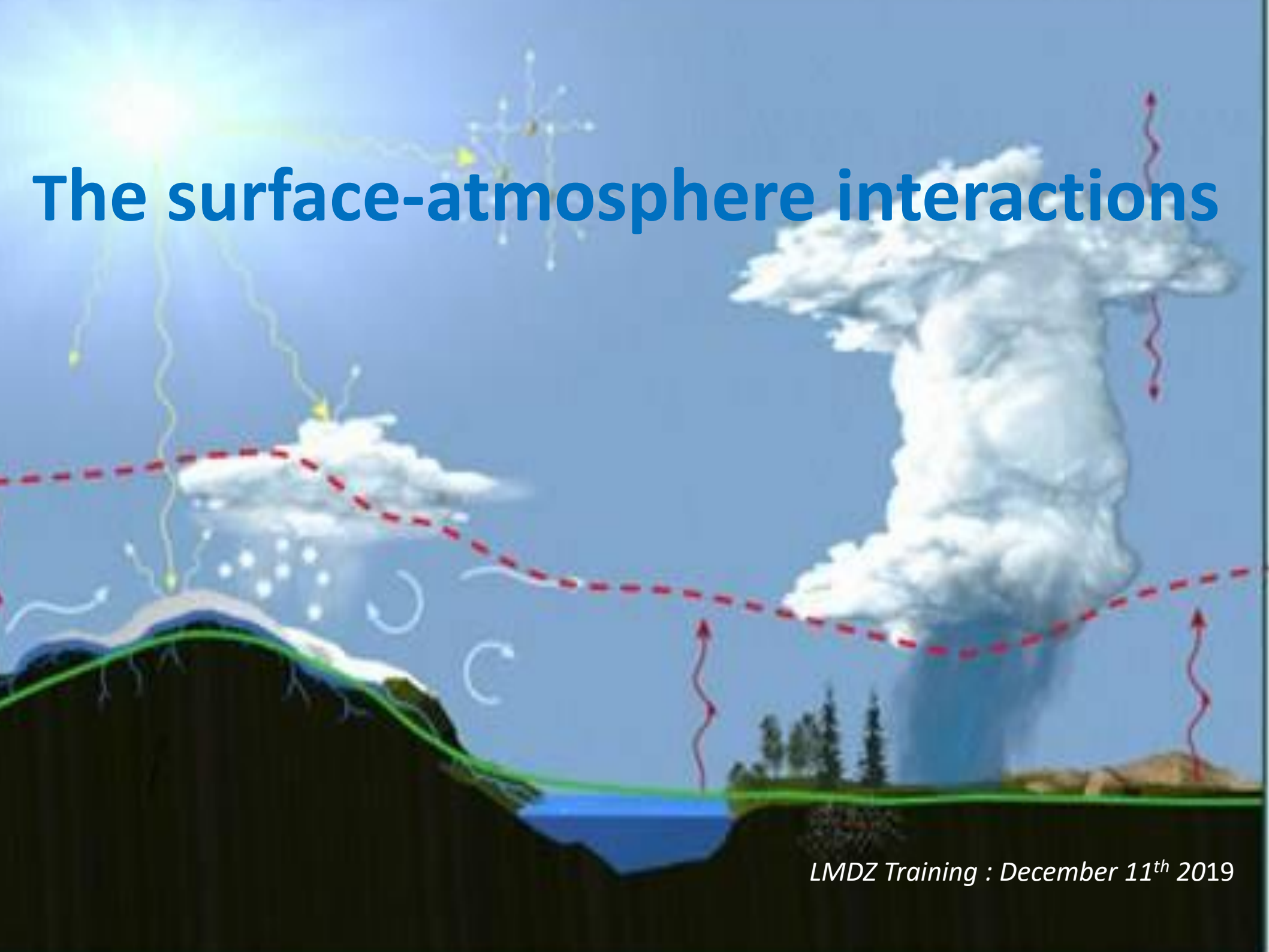
The surface-atmosphere interactions

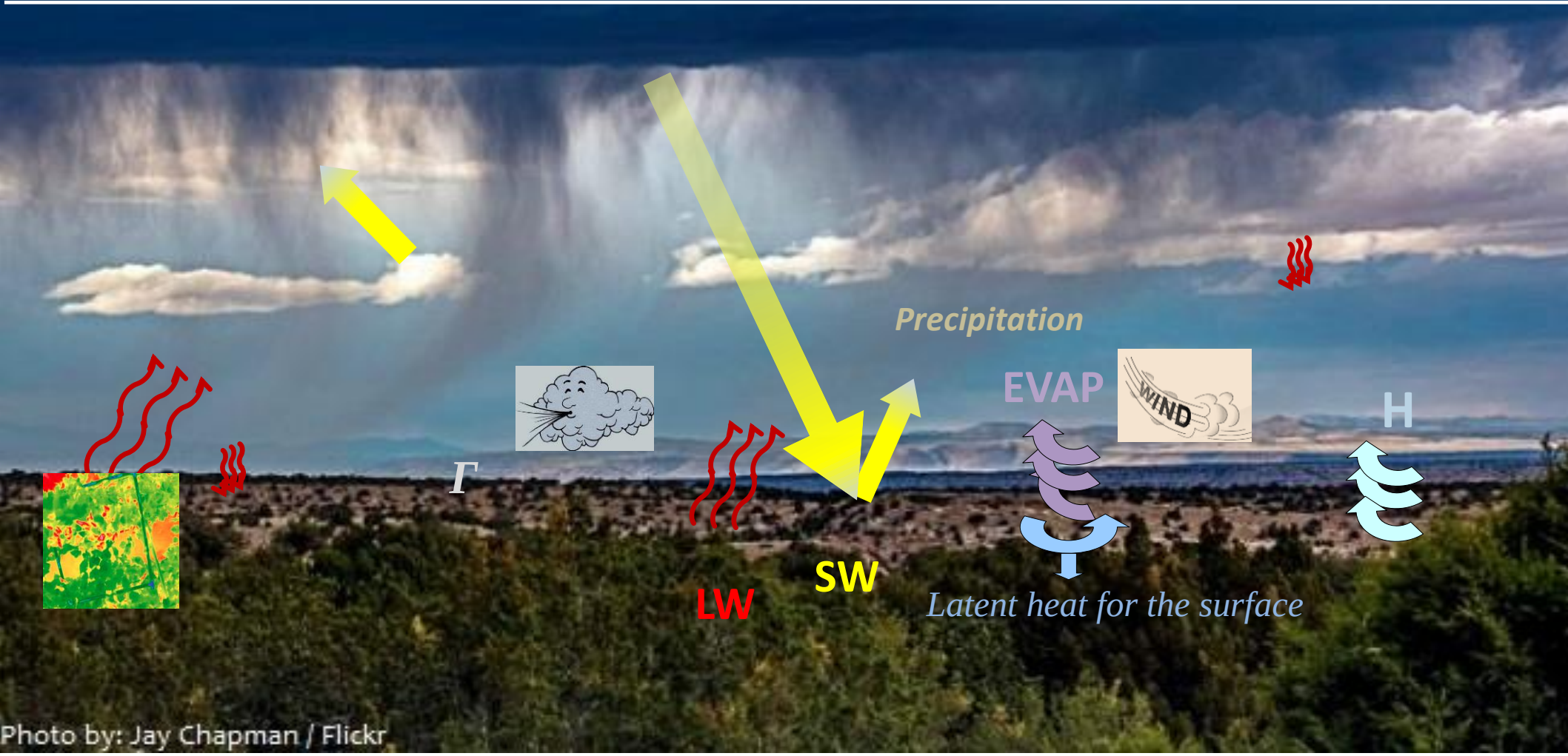


LMDZ Training : December 11th 2019

Why the surface atmosphere interactions are important ?

- To describe the climate = One needs to describe interactions between Ocean-Atmosphere-Ice caps-land surfaces under the external solar forcing and the modified atmospheric composition (CO2, aerosols, CH4 ..)
- First order Ocean-Atmosphere interactions dominate (exchange of heat, moisture ..)
- The land-surface interactions are essential for the high frequency variability of near surface meteorology, can strongly modulate the regional climates, impact the hydrological cycle…
- Partially control climate hazards: Heat waves, droughts … and their consequences e.g famines

COUPLING BETWEEN ATMOSPHERE AND SURFACE



The **atmosphere and the surface are coupled** through *turbulence* (in boundary layer) and *radiation*
(SW and LW). Currently, there is no direct influence of the surface to other parametrizations.

The surface “receive” precipitation from the atmosphere (no direct feedback).

Atmosphere-surface interactions

The **atmosphere and the surface are coupled** through *turbulence* (in the boundary layer) and *radiation* (SW and LW). Currently, there are no direct influence of the surface to other parametrizations.

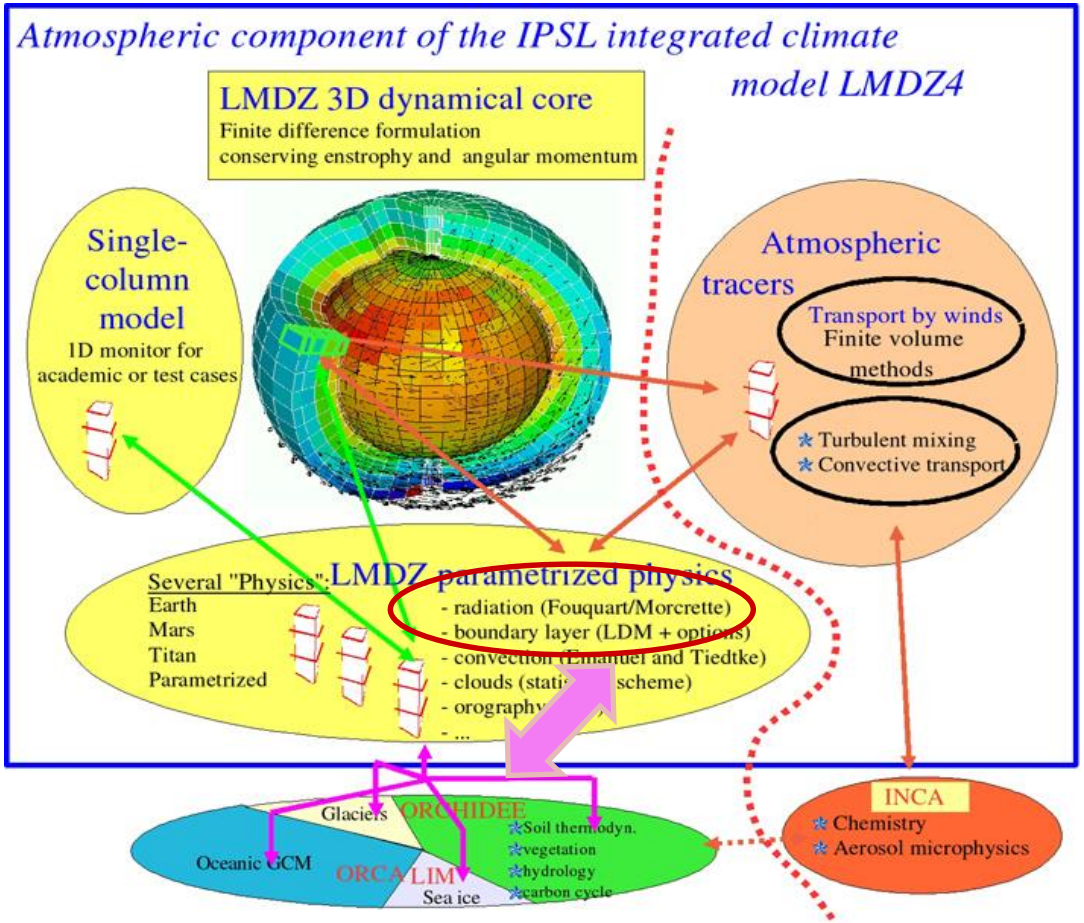
The surface “receive” precipitation from the atmosphere (no direct feedback).

The surface impacts the atmosphere via the orography (factors constant with time), roughness, albedo emissivity

In LMDZ:
Each surface grid can be decomposed in a Maximum of 4 sub-grid of different type:
land (_ter), continental ice (_lic),
open ocean (_oce) and sea_ice (_sic)

Radiation at the surface depends on mean surface properties (albedo, emissivity)

Turbulent diffusion depends on local sub-grid properties but each sub-surface sees the same atmosphere



Components of the surface energy budgets

Turbulent diffusion (pbl_surface)

Change of a variable X with the time due to the turbulent transport (continuity) :

$$\frac{\partial X}{\partial t} = -\frac{1}{\rho}\frac{\partial \Phi}{\partial z} = -\frac{\partial \Phi}{m_l}$$

m_l = mass per surface unit (kg/m^2)

X= specific humidity, momentum, moist static energy, tracers

$$\Phi = -\rho k_z \frac{\partial X}{\partial z}$$

k_z Diffusion coefficient (m^{2}s^{-1})

Φ: upward positive

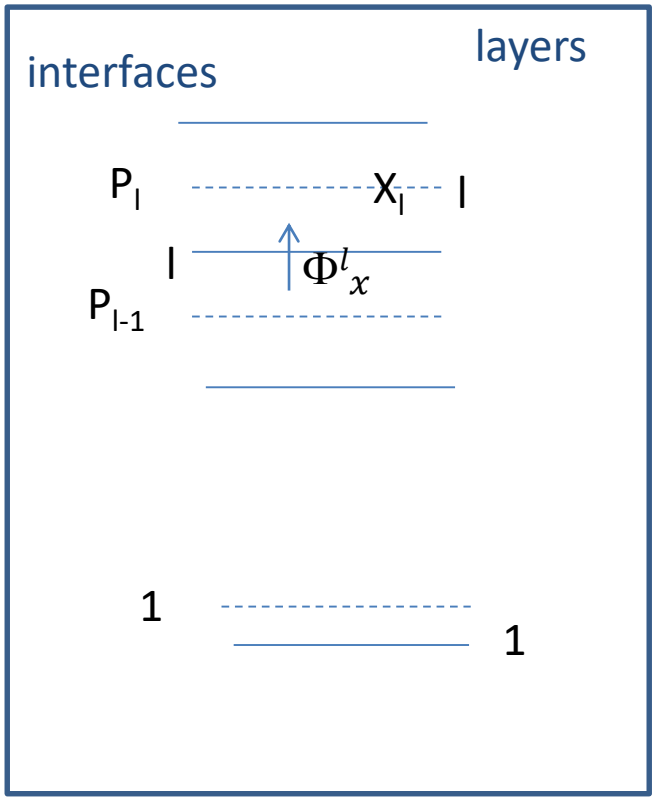
- Vertical discretization

$$\Phi^l = -K_l (X_l - X_{l-1})$$

Calculated by the turbulent scheme (controlled with iflag_pbl)

$\delta P = (P_{l-1} - P_l) = \rho g \delta z = m_l g$

$$K_l = \frac{k_z \rho^2 g}{P_{l-1} - P_l}$$



vertical discretization

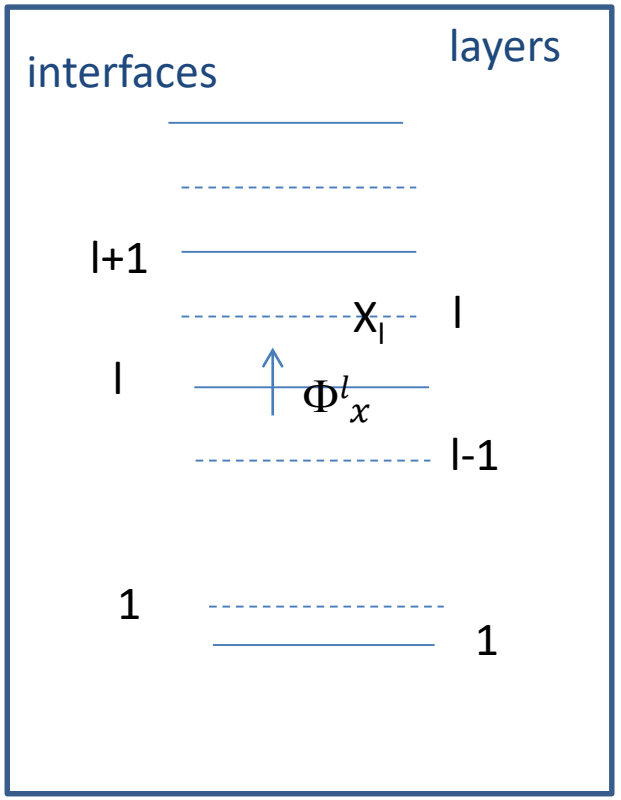
$$\frac{\partial X}{\partial t} = -\frac{\partial \Phi}{m_l} \qquad \Phi^l{}_x = -K_l\,(X_l - X_{l-1})$$

$$m_l \frac{X_l(t+\delta t) - X_l(t)}{\delta t} = \phi_l(t+\delta t) - \phi_{l+1}(t+\delta t)$$

$$m_l \frac{X_l - X_l^0}{\delta t} = \phi_l - \phi_{l+1} \quad \text{with} \quad X_l = X_l(t+\delta t),\; X_l^0 = X_l(t)$$

$$m_l \frac{X_l - X_l^0}{\delta t} = K_{l+1}(X_{l+1} - X_l) - K_l(X_l - X_{l-1})$$

$$-K_l X_{l-1} + \left(\frac{m_l}{\delta t} + K_{l+1} + K_l\right) X_l + K_{l+1} X_{l+1} = \frac{m_l}{\delta t} X_l^0$$



Tri-diagonal system that can be solved for the vector X

TOP

$$X_n = C_n^X + D_n^X X_{n-1}$$

$$C_n^X = \frac{X_n^0 \delta P_n}{\delta P_n + R_n^X}$$

$$D_n^X = \frac{R_n^X}{\delta P_n + R_n^X}$$

$2 \leq l < n$

$$X_l = C_l^X + D_l^X X_{l-1}$$

$$C_l^X = \frac{X_l^0 \delta P_l + R_{l+1}^X C_{l+1}^X}{\delta P_l + R_l^X + R_{l+1}^X (1 - D_{l+1}^X)}$$

$$D_l^X = \frac{R_l^X}{\delta P_l + R_l^X + R_{l+1}^X (1 - D_{l+1}^X)}$$

depend only on pro
layers above and th
the previous time st

$l = 1$

$$X_1 = A_1^X + B_1^X . F_1^X . \delta t$$

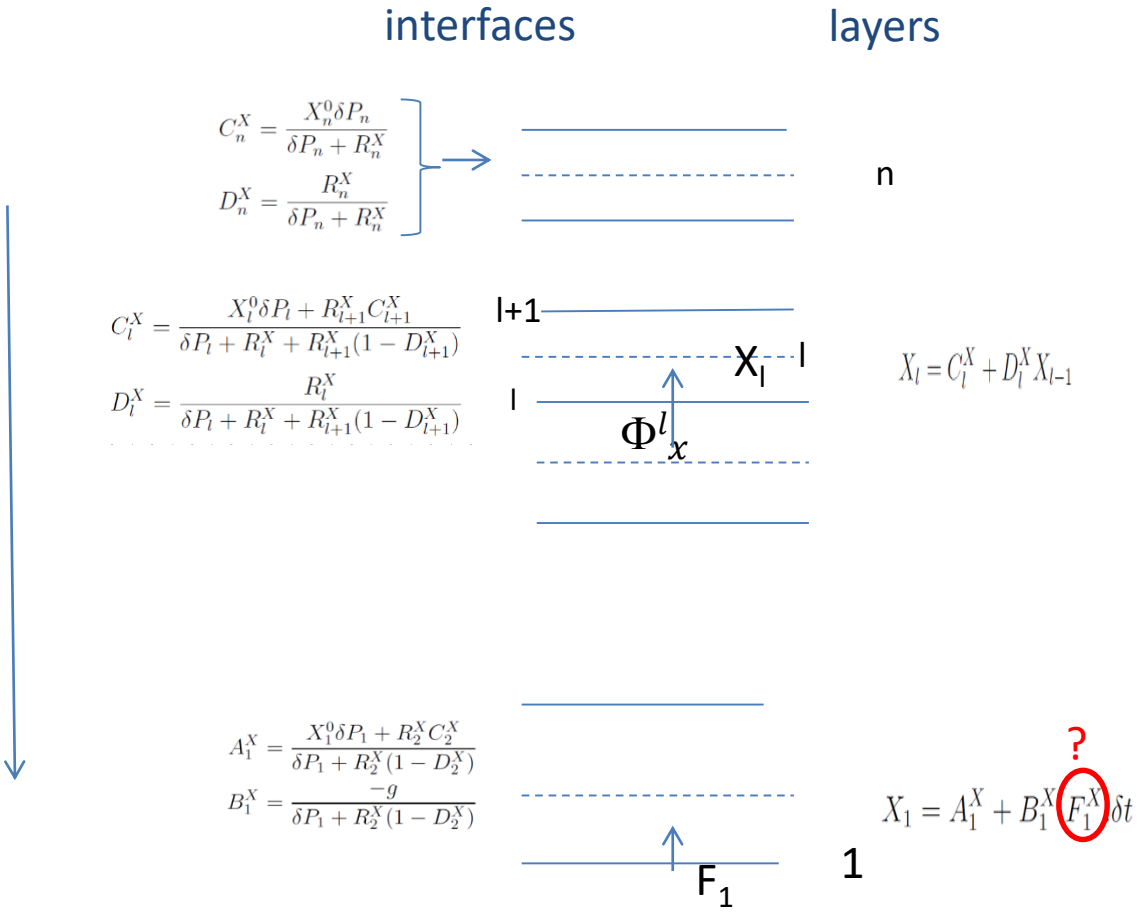
$$A_1^X = \frac{X_1^0 \delta P_1 + R_2^X C_2^X}{\delta P_1 + R_2^X (1 - D_2^X)}$$

$$B_1^X = \frac{-g}{\delta P_1 + R_2^X (1 - D_2^X)}$$

Solving the tridiagonal system

In LMDZ
climb_hq_down
climb_wind_down

downhill



X= wind, enthalpie, specific humidity, tracers

$F_1{}^x$ (flux of water mass, flux of heat, flux of momentum) is either prescribed or computed for each sub-surface

Once $F_1{}^x$ is known, the X_i can be computed from the first layer to the top of the PBL

Coupling with the surface : Compute F_x^1

A_1 and B_1 rely on the vertical diffusion scheme

$$F^1_X = \rho\, C_d^X |V| (X_1 - X_s) = \frac{X_1 - A_1}{B_1 \delta t}$$

$$X_1 = A_1^X + B_1^X . F_1^X . \delta t$$

Value of X_1 if F_1=0

Sensitivity of X_1 to F_1

C_d^X drag coefficient (Monin Obukhov, constant flux in the surface layer) routine cdrag.F90 depends on

- roughness lenghts (gustiness, vegetation), ~~orography~~
- Richardson number (boundary layer stability)
- Formulation depends on the sub-surface type

Implicit coupling: Surface conditions (X_s) should be evaluated at the same time as temperature, humidity in the atmosphere (X_1)

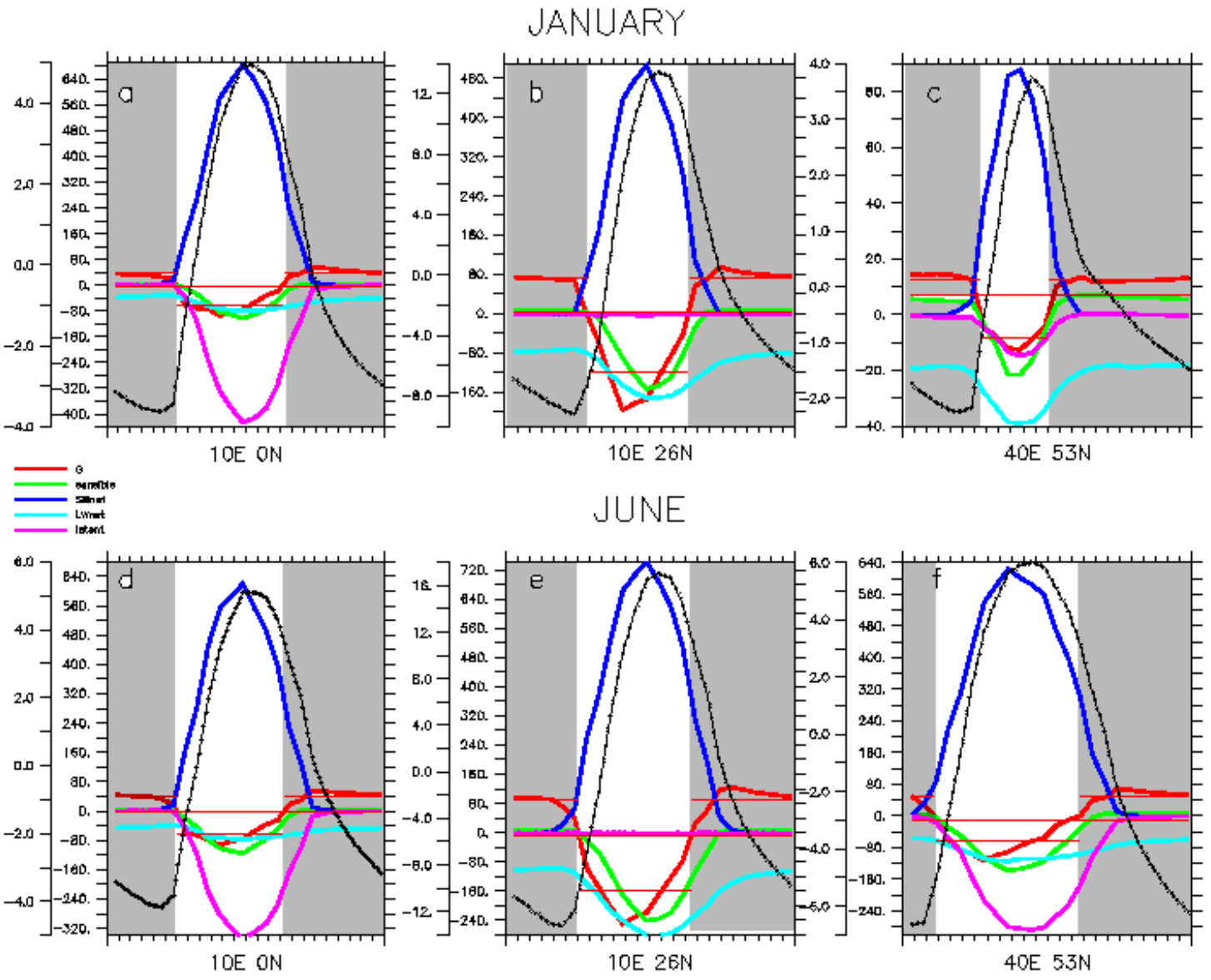
Once $\mathbf{X_s}$ is known , X_1 and F^1_X are known

$$H = \overline{\rho} C_p \overline{w'T'},$$

$$LE = L\overline{\rho}\overline{w'q'},$$

Case of the continental surface

$$SW_{net} + LW_{net} + H + L + \Phi_0 = 0$$



SIRTA atmospheric observatory 48.7° N – 2.2° E
Dynamics and turbulence:
X-band radar
T-RH profile
Doppler lidar, sodar, UHF radar
Sonic anemometers
Clouds and aerosols:
UV-VIS-NIR lidars (depol, retro.)
Sun-photometer
Solar and infrared radiometers
Sky-cloud imager
FMCW cloud radar
Microwave radiometer
Aerosol mass spectrometer
LMD
Routine observations

Case of the continental surface

$$SW_{net} + LW_{net} + H + L + \Phi_0 = 0$$

$$SW_{net} + LW_d - \varepsilon\sigma T_s^4 + H + L + \Phi_0 = 0$$

depends on Ts

$$H = \beta\rho VC_d\,(q_1 - q_s(T_s))$$

$$L = \rho VC_d\,(T_1 - T_s)$$

- Heat conduction in the soil: diffusion equation :

$$\begin{cases} \Phi_T = -\lambda \dfrac{\partial T}{\partial z} \\ \dfrac{\partial T}{\partial t} = -\dfrac{1}{C} \dfrac{\partial \Phi_T}{\partial z} \end{cases}$$

λ= thermal diffusivity
C = thermal capacity

Boundary conditions:

- ✓ bottom : $\Phi = 0$
- ✓ top: Continuity of the fluxes and the temperature between sub-surface and atmosphere

Case of the continental surface

- Heat conduction in the soil: diffusion equation :

$$\begin{cases} \Phi_T = -\lambda \dfrac{\partial T}{\partial z} \\ \dfrac{\partial T}{\partial t} = -\dfrac{1}{C}\dfrac{\partial \Phi_T}{\partial z} \end{cases}$$

+

- bottom boundary condition $\Phi = 0$

→ solve for the temperature into the soil (tridiagonal system)

- Top boundary condition:

Continuity of the fluxes and the temperature between sub-surface and atmosphere

$$C' * \frac{T_s^t - T_s^0}{\delta t} = G' * + SW_{net} + LW_d + \sum F^{\downarrow}(T_S^t) - \varepsilon\sigma(T_S^t)^4$$

Solved using the sensitivity of the flux to the surface temperature to calculate the flux at the new time-step

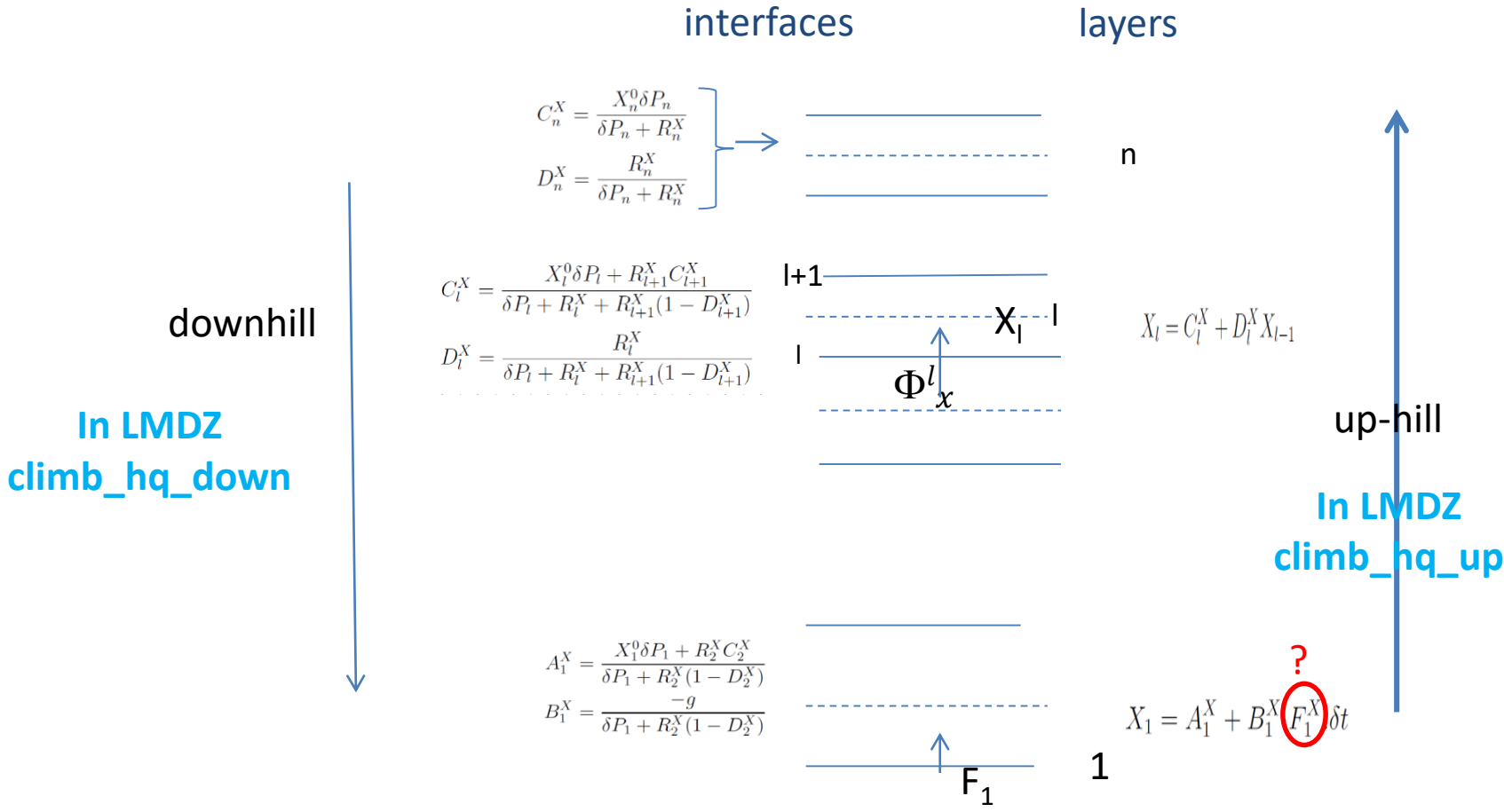
$$F_{s,H}^t = sensfl_{old} - sensfl_{sns}(T_s^t - T_s^{t-\delta t})$$

$$\sigma * T_s^{t-\delta t\,4} - 4\varepsilon\sigma T_s^{t-\delta t\,3}(T_s^t - T_s^{t-\delta t})$$

$$T_S^t = f(SW_{net} + LW_d, T_s^0, F_s^0, A_1, B_1)$$

fsoil called by surf_land_bucket_mod in LMDZ
enerbil and thermosoil when coupled w.Orchidee

Solving the tridiagonal system

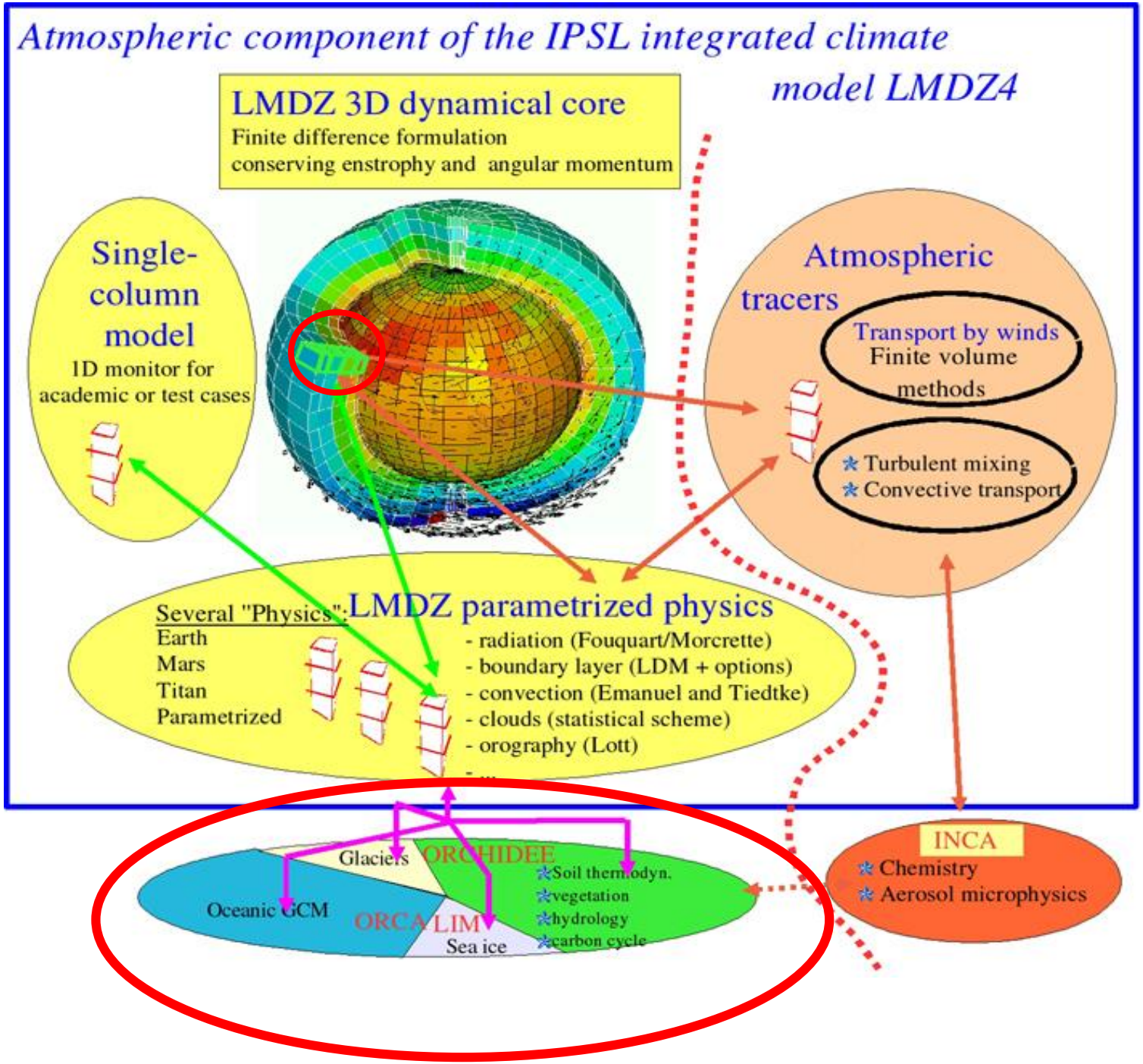


X= wind, enthalpie, specific humidity, tracers

F_1^x (flux of water mass, flux of heat, flux of momentum) is either prescribed or computed for each sub-surface $F^1_X = \rho C_d^x |V| (X_1 - X_s) = \frac{X_1 - A_1}{B_1 \delta t}$

$$F^t_{s,H} = sensfl_{old} - sensfl_{sns}(T^t_s - T^{t-\delta t}_s)$$

Once F_1^x is known, the X_i can be computed from the first layer to the top of the PBL

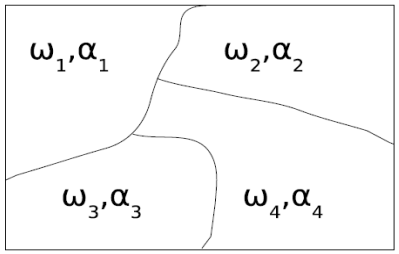

Atmospheric component of the IPSL integrated climate model LMDZ4
LMDZ 3D dynamical core
Finite difference formulation
conserving enstrophy and angular momentum
Single-column model
1D monitor for academic or test cases
Atmospheric tracers
Transport by winds
Finite volume methods
Turbulent mixing
Convective transport
Several "Physics":
Earth
Mars
Titan
Parametrized
LMDZ parametrized physics
- radiation (Fouquart/Morcrette)
- boundary layer (LDM + options)
- convection (Emanuel and Tiedtke)
- clouds (statistical scheme)
- orography (Lott)
Glaciers
ORCHIDEE
Soil thermodyn.
vegetation
hydrology
carbon cycle
Oceanic GCM
ORCA LIM
Sea ice
INCA
Chemistry
Aerosol microphysics

Coupling between atmospheric column(s) and sub-surfaces

Each grid cell is divided into several sub-areas or ``sub-surfaces" of fractions ω_i

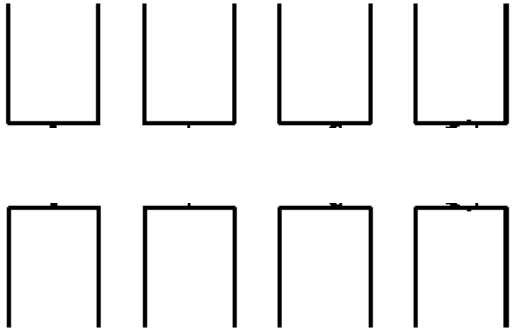
Sub-surfaces

$$\sum_i \omega_i = 1$$



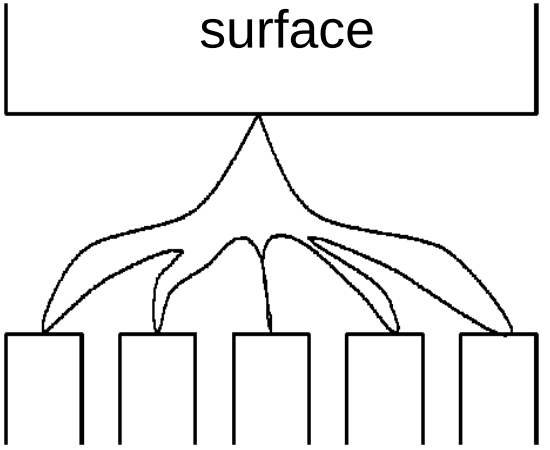
Turbulent flux

One PBL over **each** sub-surface



Radiative flux

One column **covers all** the sub-surface



Each sub surface has to compute F_1 using variables X_1, A_1 and B_1

The boundary layer tendencies in the atmosphere are mixed between sub-columns (equivalent of averaging the surface flux)

Derivation of local sub-surface **net solar radiation** from grid average net solar radiation

The grid average net flux *Ψs* at surface has been computed for each grid point by the radiative cod

We want (1) to conserve energy and (2) to take into account the value of the local albedo αi of the s

We compute the downward SW radiation as

with the mean albedo $\alpha = \sum_i \omega_i \alpha_i$ $F_{\downarrow}^{s} = \frac{\Psi_s}{(1-\alpha)}$

For each sub-surface i, the absorbed solar radiation reads: $\psi_i^s = (1-\alpha_i) F_{\downarrow}^{s}$

$$\psi_i^s = \frac{(1-\alpha_i)}{(1-\alpha)} \bar{\Psi}^s$$

One may verified that this procedure ensure energy conservation, $\sum_i \omega_i \psi_i^s = \Psi_s$

Derivation of local sub-surface **net longwave radiation** from grid average net longwave radiation

The net longwave (LW) radiation $\bar{\Psi}^L$ has been computed by the radiative code for each grid cell. How to split it depending on the sub-surfaces local properties and ensuring energy conservation?

If the downward longwave flux $F_\downarrow$ is uniform within each grid, the net LW flux for a sub-surface i may be written as:

$$\psi_i^L(T_i) = \epsilon_i \left(F_\downarrow - \sigma T_i^4\right) \tag{1}$$

where T_i is the surface temperature of sub-surface i and ϵ_i its emissivity. A linearization around the mean temperature $\bar{T}$ gives:

$$\psi_i^L(T_i) \approx \epsilon_i \left(F_\downarrow - \sigma \bar{T}^4\right) - 4\epsilon_i \sigma \bar{T}^3 (T_i - \bar{T}) \tag{2}$$

To conserve the energy, the following relationship must be true:

$$\sum_i \omega_i \psi_i^L = \bar{\Psi}^L \tag{3}$$

Using Eq. 2 gives

$$\sum_i \omega_i \psi_i^L = \bar{\epsilon} \left(F_\downarrow - \sigma \bar{T}^4\right) - 4\sigma \bar{T}^3 \sum_i \omega_i \epsilon_i \left(T_i - \bar{T}\right) \tag{4}$$

where $\bar{\epsilon} = \sum_i \omega_i \epsilon_i$ is the mean emissity.

Derivation of local sub-surface **net longwave radiation** from grid average net longwave radiation

$$\sum_i \omega_i \psi_i^L = \bar{\epsilon}\left(F_\downarrow - \sigma \bar{T}^4\right) - 4\sigma \bar{T}^3 \sum_i \omega_i \epsilon_i \left(T_i - \bar{T}\right) \tag{4}$$

where $\bar{\epsilon} = \sum_i \omega_i \epsilon_i$ is the mean emissity. The second term on the right hand side is zero if

$$\bar{T} = \frac{\sum_i \omega_i \epsilon_i T_i}{\bar{\epsilon}} \tag{5}$$

To ensure energy conservation, we need in addition to verify:

$$\bar{\Psi}^L = \bar{\epsilon}\left(F_\downarrow - \sigma \bar{T}^4\right) \tag{6}$$

Which is consistent with the definition of the net LW flux at the surface. We rewrite now Eq. 2 as:

$$\psi_i^L(T_i) \approx \frac{\epsilon_i}{\bar{\epsilon}} \bar{\Psi}^L - 4\epsilon_i \sigma \bar{T}^3 (T_i - \bar{T}) \tag{7}$$

Due to radiative code limitation, in LMDZ, we always must have $\varepsilon_i = 1$

Call tree

In subroutine PHYSIQ
loop over time steps
CALL change_srf_frac : Update fraction of the sub-surfaces (pctsrf)
....
CALL pbl_surface Main subroutine for the interface with surface
Calculate net radiation at sub-surface
Loop over the sub-surfaces nsrf
Compress variables (Consider only one surface type and only the points for which the fraction for this sub-surface in not zero)
CALL cdrag: coefficients for turbulent diffusion at surface (cdragh and cdragm)
CALL coef_diff_turb: coef. turbulent dif. in the atmosphere (ycoefm et ycoefm.)
CALL climb_hq_down downhill for enthalpy H and humidity Q
CALL climb_wind_down downhill for wind (U and V)
CALL surface models for the various surface types: **surf_land, surf_landice, surf_ocean or surf_seaice**. **Each surface model computes:**
• evaporation, latent heat flux, sensible heat flux
• surface temperature, albedo (emissivity), roughness lengths
CALL climb_hq_up : compute new values of enthalpy H and humidity Q
CALL climb_wind_up : compute new values of wind (U and V)
Uncompress variables : (some variables are per unit of sub-surface fraction, some are per unit of grid surface fraction)
Cumulate in global variables after weighting by sub-surface fractions
Surface diagnostics : (T, q, wind are evaluated at a reference level (2m) owing to an interpolation scheme based on the MO laws).
End Loop over the sub-surfaces
Calculate the mean values over all sub-surfaces for some variables
End pbl-surface

About the surface models: land, ocean, sea-ice , land-ice

For land : simplified land surface model, hydrology= bucket or “beta clim”,
constant thermal inertia (soil /snow) , albedo and rugosity from a file.
(surf_land_bucket: beta=min(2*qsol/max_eau_sol,1))

or Soil-vegetation-atmosphere transfer (SVAT) model (ORCHIDEE)

For ocean: Forced, fully coupled (with NEMO), coupled with a slab-ocean

For sea ice depends on the coupling with the ocean (forced, coupled, slab)

For land-ice : snow properties calculated with sisvat if ok_snow= T.
otherwise simplified (as for land + simplified snow prop., rugo, albedo)

About the values interpolated at a reference level near the surface (e.g. 2m)

Principle: Constant flux in the surface layer and similarity laws: Non dimensional vertical gradient of horizontal wind, potential temperature, specific humidity are assumed to be universal function of a stability parameter z/L (L= Monin-Obukhov law) or of the Richardson Number (Louis 82).

$$\frac{kz}{u_*}\frac{\partial u}{\partial z} = \phi_M\left(\frac{z}{L}\right) \quad ; \quad \frac{kz}{\theta_*}\frac{\partial \theta}{\partial z} = \phi_H\left(\frac{z}{L}\right) \quad ; \quad \frac{kz}{q_*}\frac{\partial q}{\partial z} = \phi_\varphi\left(\frac{z}{L}\right)$$

integrating these equations gives

$$\frac{ku}{u_*} = \ln\left(\frac{z}{z_0}\right) - \psi_M\left(\frac{z}{L}\right) + \psi_M\left(\frac{z_0}{L}\right)$$

$$\frac{k(\theta-\theta_0)}{\theta_*} = \ln\left(\frac{z}{z_0}\right) - \psi_H\left(\frac{z}{L}\right) + \psi_H\left(\frac{z_{0H}}{L}\right)$$

Louis shows that one can use the Richardson (bulk) number instead of the monin obukhov length

$$\frac{ku}{u_*} = \frac{\ln(z/z_0)}{F_M^{1/2}\left(R_i^b, \frac{z}{z_0}\right)}$$

$$\frac{k(\theta-\theta_0)}{\theta_*} = \frac{\ln\left(\frac{z}{z_0}\right)}{F_H\left(R_i^b, \frac{z}{z_0}\right)} F_M^{1/2}\left(R_i^b, \frac{z}{z_0}\right)$$

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$$\frac{k\,u}{u_*} = \frac{\ln(z/z_0)}{F_M^{1/2}(R_i^b, \frac{z}{z_0})}$$

$$\text{II} \quad \frac{k(\theta-\theta_0)}{\theta_*} = \frac{\ln\left(\frac{z}{z_0}\right)}{F_H(R_i^b, \frac{z}{z_0})} F_M^{1/2}(R_i^b, \frac{z}{z_0})$$

if one writes II for 2 levels (first atm. level and reference level)

$$\frac{\theta_{ref}-\theta_0}{\theta_1+\theta_0} = \frac{\ln\left(\frac{z_{ref}}{z_0}\right)}{\ln\left(\frac{z_1}{z_0}\right)} \cdot \frac{F_M^{1/2}(R_i^{ref}, \frac{z_{ref}}{z_0})}{F_M^{1/2}(R_i^1, \frac{z_1}{z_0})} \cdot \frac{F_H(R_i^1, \frac{z_1}{z_0})}{F_H(R_i^{ref}, \frac{z_{ref}}{z_0})}$$

$\theta_0 = T_s$

→ evaluate $\theta_{ref} = \varphi(\theta_0, \theta_1, \text{stability}, z_1, z_{ref}, R_i^1, R_i^{ref})$

$$\frac{q_{ref}+q_0}{q_1-q_0} = \frac{\ln\left(\frac{z_{ref}}{z_0}\right)}{\ln\left(\frac{z_1}{z_0}\right)} \cdot \frac{F_M^{1/2}(R_i^{ref}, \frac{z_{ref}}{z_0})}{F_M^{1/2}(R_i^1, \frac{z_0}{z_0})} \cdot \frac{F_H(R_i^1, \frac{z_1}{z_0})}{F_H(R_i^{ref}, \frac{z_{ref}}{z_0})}$$

Some properties –

stability function in the stable case : $\psi_H\left(\frac{z}{L}\right) = -\frac{5}{L}$

$(L > 0)$

$$\frac{\theta - \theta_0}{\theta_*} = \frac{1}{k}\left[\ln\frac{z}{z_0} + 5\frac{z}{L}\right]$$

at fixed L, $\theta_* = cst$ and does not depend on z.

$$\Rightarrow \frac{d\,\delta\theta}{\delta z} = \theta_* = \text{cste} = \text{function monotone of } z$$

For a stability: $\theta_s \leq \theta \leq \theta_1$

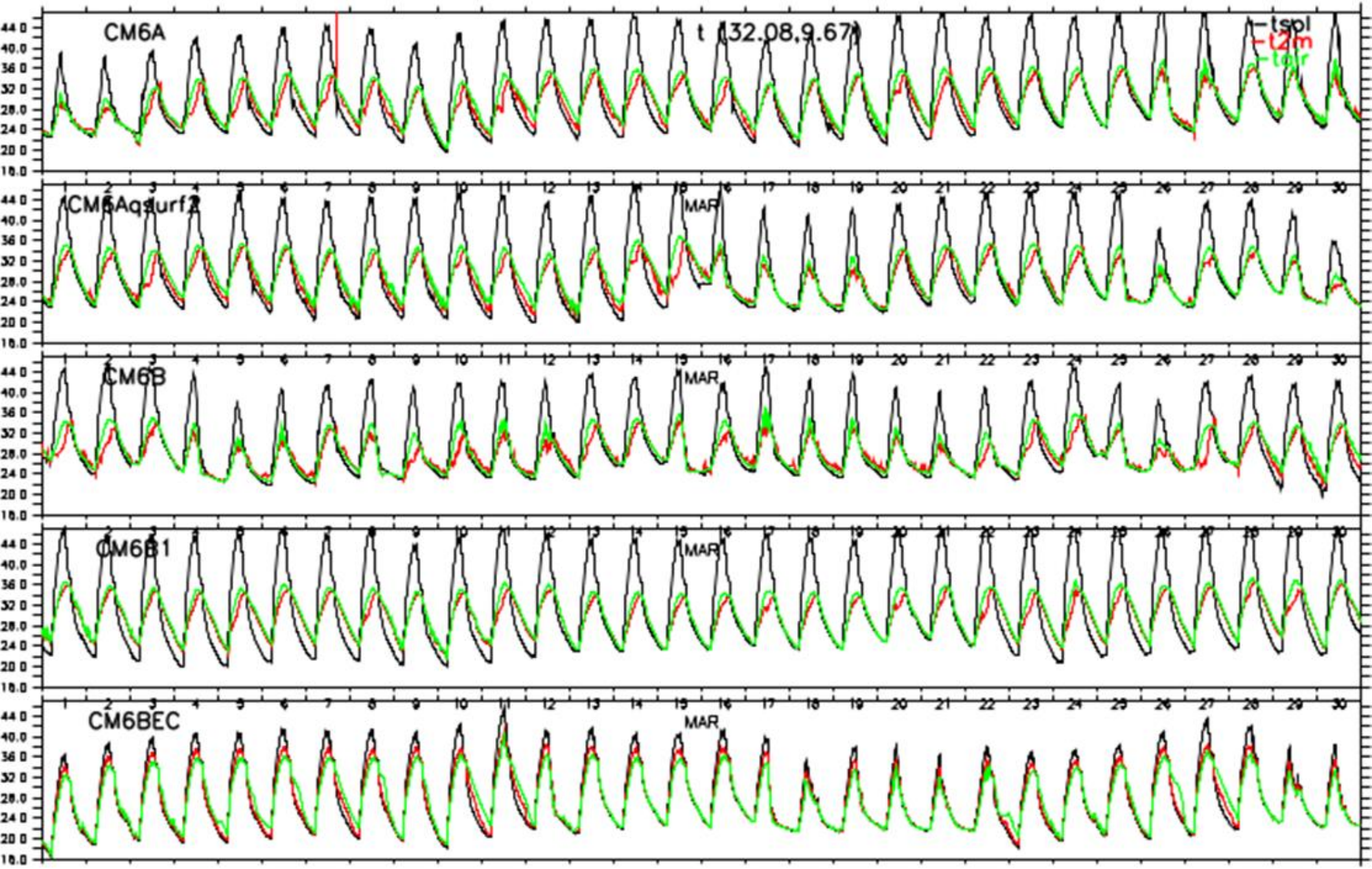
In the unstable case : $\frac{\theta - \theta_0}{\theta_*} = \ln\left(\frac{z}{z_0}\right) + 2\ln\left[\left(1 + \left(1 - 16\frac{z}{L}\right)^{\frac{1}{2}}\right)\right]$

$L < 0$

$\frac{d\,\delta\theta}{\delta z}$ not always monotone –

$$F\left(\frac{z}{L}\right) : -\frac{16}{L} \cdot \frac{1}{\left(1 - 16\frac{z}{L}\right)^{1/2}\left[1 + \left(1 - 16\frac{z}{L}\right)^{1/2}\right]}$$

$$F'\left(\frac{z}{L}\right) < 0 .$$


CM6A
t (32.08,9.67)
tspl
t2m
tsr
CM6Aqsurf2
CM6B
CM6B1
CM6BEC
MAR

- Technical note **:** Description of the interface with the surface and the computation of the turbulent diffusion in LMDZ (J.L.Dufresne) Thèse F. Hourdin 1993 (section 3.3.3 and annexes)
- Wang F., F. Cheruy, J.L. Dufresne, 2016: The improvement of soil thermodynamics and its effects on land surface meteorology in the IPSL climate model. Geosci. Model Dev., 9, 363–381, 2016 www.geosci-model-dev.net/9/363/2016/
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Solving the tridiagonal system

$$\left(\frac{m_l}{\delta t} + K_{l+1} + K_l\right) X_l = \frac{m_l}{\delta t} X_l^0 + K_{l+1} X_{l+1} + K_l X_{l-1}$$

which may be written as:

$$\left(\delta P_l + R_{l+1}^X + R_l^X\right) X_l = \delta P_l \; X_l^0 + R_{l+1}^X \; X_{l+1} + R_l^X \; X_{l-1} \qquad (2 \leq l < n)$$

$$\text{with } R_l^X = g\delta t K_l$$

At the top (l=n, $\Phi_n = 0$)

$$\left(\delta P_n + R_n^X\right) X_n = \delta P_n \; X_n^0 + R_n^X \; X_{n-1}$$

At the bottom: (l=1) : $m_1 \frac{x_1 - x^0_1}{\delta t} = \Phi^1_x - \Phi^2_x$

$$m_1 \frac{X_1 - X_1^0}{\delta t} = K_2(X_2 - X_1) - F_1^X$$

$$\left(\delta P_1 + R_1^X\right) X_1 = \delta P_1 \; X_1^0 + R_2^X \; X_2 - g\delta t F_1^X$$

With F_1^X : flux of X at the bottom of the first layer (i.e. between the surface and the atmosphere), positive downward.

$$K_l = \frac{k_z \rho^2 g}{P_{l-1} - P_l}$$